

Beyond the mean-field limit for the McKean-Vlasov particle system: Uniform in time estimates for the cumulants

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$$dX_t = b(t, X_t)dt + \sigma(t, X_t)dB_t,$$

with

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Today,  $\sigma \equiv I_d$  (the identity matrix) to simplify.

## A PDE interpretation of the McKean-Vlasov equation

McKean-Vlasov SDE

$$X_t = X_0 + \int_0^t b(X_s, \mathcal{L}(X_s)) ds + B_t, \quad t \geq 0$$

with  $\mathcal{L}(X_0) = \mu_0$  for some distribution  $\mu_0$ .

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Starting from  $\mu_0$ , the flow of marginal laws  $(m(t, \mu_0) := \mathcal{L}(X_t))_{t \geq 0}$  satisfies a nonlinear Fokker-Planck equation:

$$\begin{aligned} \partial_t m(t, \mu) &= \frac{1}{2} \Delta m(t, \mu) - \operatorname{div} [m(t, \mu) b(\cdot, m(t, \mu))], \quad t \geq 0, \\ m(0, \mu) &= \mu, \end{aligned}$$

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This is a *PDE interpretation* of the McKean-Vlasov SDE.

## The particle system: first-order, MF interaction and noise

Key question: deriving the previous McKean-Vlasov SDE from a microscopic system. Given  $N$  particles (agents), can we obtain the previous SDE as a limit ? In which sense ?

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For particles in the torus  $\mathbb{T}^d$ :

$$\left\{ \begin{array}{l} Y_t^{i,N} = Y_0^{i,N} + \int_0^t b(Y_s^{i,N}, \mu_s^N) ds + B_t^i, \\ \mu_s^N := \frac{1}{N} \sum_{i=1}^N \delta_{Y_s^{i,N}}, \end{array} \right. \quad t \geq 0, \quad 1 \leq i \leq N,$$

where

- ☞  $(Y^{i,N})_{i=1}^N$  are the positions;
- ☞  $\mu_s^N$  empirical measure at time  $s$ ;
- ☞  $b : \mathbb{T}^d \times \mathcal{P}(\mathbb{T}^d) \rightarrow \mathbb{R}^d$  is an interaction potential and the **mean-field** scaling is considered.

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Additionally: smooth  $b$ ,  $(Y_0^{i,N})_{1 \leq i \leq N} \sim \mu_0^{\otimes N}$  (i.i.d. initial distributions). The goal is to **relate this particle system to the McKean-Vlasov SDE**.

## Assumptions on $b$

H-stable potential (Carrillo, Gvalani, Pavliotis, Schlichting 2020)

$$b(x, m) = -\kappa \int_{\mathbb{T}^d} \nabla W(x - y) m(dy), \quad x \in \mathbb{T}^d, m \in \mathcal{P}(\mathbb{T}^d)$$

for  $\kappa > 0$  (equal to 1 in what follows for simplicity) and  $W$  smooth, coordinate-wise even:

$$W(x_1, \dots, -x_i, \dots, x_d) = W(x_1, \dots, x_i, \dots, x_d), \quad (x_1, \dots, x_d) \in \mathbb{T}^d.$$

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Consequence: the Lebesgue measure on  $\mathbb{T}^d$ ,  $\text{Leb}_{\mathbb{T}^d}$  is

- ✓ the unique invariant measure for the McKean-Vlasov equation;
- ✓ exponentially stable, i.e. there exists  $C, \lambda > 0$  constants s.t. for all  $t \geq 0$ ,

$$\|m(t, \mu) - \text{Leb}_{\mathbb{T}^d}\|_{TV} \leq C e^{-\lambda t}.$$

## Statistical description and mean-field limit

$N \gg 1$ ,  $F_N(t, x_1, \dots, x_N)$  probability density on the  $N$  torus  $(\mathbb{T}^d)^N$  at time  $t \geq 0$   
and  $F_N^1$  first marginal

$$F_N^1(t, z) = \int_{(\mathbb{T}^d)^{N-1}} F_N(t, z, z_2, \dots, z_N) dz_2 \dots dz_N.$$



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When  $N \rightarrow \infty$ , in view of **Boltzmann chaos assumption** one wants to **neglect the correlations** and to obtain, in the limit  $N \rightarrow \infty$ , that  $F_N^1$  behaves like the solution of the McKean-Vlasov SDE.

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Can we justify this convergence ?

## BBGKY, MV version: a formal expansion

$F^N$  solves a forward Kolmogorov equation, where  $\bar{x} = (x_1, \dots, x_N) \in (\mathbb{T}^d)^N$

$$\partial_t F^N(\bar{x}) = \frac{1}{2} \Delta F^N(\bar{x}) + \sum_{i=1}^N \operatorname{div}_{x_i} \left( F^N(\bar{x}) \frac{1}{N} \sum_{j=1}^N \nabla W(x_j - x_i) \right).$$

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Integrating with respect to  $x_2, \dots, x_N$ , writing  $F_N^k$  for the  $k$ -th marginal, and using the coordinate-wise symmetry:

$$\partial_t F_N^1(x) = \frac{1}{2} \Delta F_N^1(x) + \operatorname{div}_x \left( \int_{\mathbb{T}^d} \nabla W(y - x) F_N^2(x, y) dy \right).$$

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Key point: at time 0,  $F_N(0, \cdot) = \mu_0^{\otimes N}(\cdot)$  so in particular

$$F_N^2(0, x, y) = F_N^1(0, x) F_N^1(0, y).$$

Not true anymore at time  $t > 0$  ! But one expects

$F_N^2(x, y) = F_N^1(x) F_N^1(y) + G_N^2(x, y)$  where  $G_N^2$  is the two-particles correlation function with  $G_N^2 \rightarrow 0$  as  $N \rightarrow \infty$ .

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Then  $\rightarrow$  PDE version of the McKean-Vlasov equation on  $\mathbb{T}^d$

$$\partial_t f(x) = \frac{1}{2} \Delta f(x) - \operatorname{div}_x \left( b(x, f) f(x) \right).$$

## Propagation of chaos I: qualitative convergence

$G_N^2 \rightarrow 0$  as  $N \rightarrow \infty$  ? Are  $Y^{1,N}$  and  $Y^{2,N}$  “independent” at all time ? In what sense ?

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In general, expect that on  $[0, T]$ ,  $T > 0$ , and for any fixed  $k \in \{1, \dots, N\}$ ,

$$(Y^{1,N}, \dots, Y^{k,N}) \Longrightarrow (X^1, \dots, X^k),$$

where  $(X^i)_i$  are i.i.d. copies of solutions to the MVSDE, weakly in  $C([0, T], (\mathbb{T}^d)^k)$ . This gives

- \* convergence towards the limit equation
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- ▶ For  $b$  Lipschitz (w.r.t. the topology of  $\mathbb{T}^d \times \mathcal{P}(\mathbb{T}^d)$ ), **Sznitman's coupling**, see Sznitman (1991), Lacker (2018)...
  - ▶ Other approaches: tightness of  $(\mathcal{L}(\mu_t^N) \in \mathcal{P}(\mathcal{P}(\mathbb{T}^d)))_{0 \leq t \leq T}$ . Then  $(\mathcal{L}(\mu_t^N))_{0 \leq t \leq T}$  converges weakly to  $\delta_{\mathcal{L}(X_t)_{0 \leq t \leq T}}$ .

## Propagation of chaos II: strong errors

Two types of results, quantifying **strong** and **weak** errors. **Second aspect to quantify: uniformity in time ?** Strong errors: convergence in some Wasserstein norm, e.g.

$$\sup_{t \geq 0} W_1 \left( F_N^k(t, \cdot), m(t, \mu_0)^{\otimes k} \right) = O\left( \frac{1}{N^{\frac{1}{2}}} \right).$$

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Recent uniform in times results:

- ▶ Malrieu 2001 ( $W$  convex);
- ▶ Durmus-Eberle-Guillin-Zimmer 2020 for small interaction;
- ▶ Guillin-Le Bris-Monmarché 2021 for more singular interactions (allowing to treat the Biot-Savart kernel).

Jabin-Wang (2018): non-uniform in time estimates for singular interaction, starting point of several papers.

Other approach to strong error: central limit theorem (Sznitman, Méléard...).

## Propagation of chaos III: weak errors

Focus on the statistical behavior of  $\mu_t^N$ . Goal: deriving rates of convergence (in  $t$  and  $N$ ) for

$$\mathbb{E} \left[ \left| \Phi(\mu_t^N) - \Phi(m(t, \mu_0)) \right| \right],$$

where  $\Phi : \mathcal{P}(\mathbb{T}^d) \rightarrow \mathbb{R}$  is a test function. Typically  $\Phi$  is

- \* polynomial: Mischler-Mouhot-Wennberg 2015;
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For the torus case, recent results of Delarue-Tse (2021): under regularity assumptions on  $b$  and  $\Phi$ , there exists  $C > 0$  such that for all  $\mu_0 \in \mathcal{P}(\mathbb{T}^d)$ ,

$$\sup_{t \geq 0} \mathbb{E} \left[ \left| \Phi(\mu_t^N) - \Phi(m(t, \mu_0)) \right| \right] \leq \frac{C}{N}.$$

## Back to the marginals

Recall:

$$\partial_t F_N^1(x) = \frac{1}{2} \Delta F_N^1(x) + \operatorname{div}_x \left( \int_{\mathbb{T}^d} \nabla W(y-x) F_N^2(x,y) dy \right).$$

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Writing  $F_N^2(x,y) = F_N^1(x)F_N^1(y) + G_N^2(x,y)$ , the previous results show  $G_N^2 = O(\frac{1}{N})$  in some weak sense  $\rightarrow$  McKean-Vlasov equation.

## Beyond mean-fields (Bogolyubov corrections?)

What if we keep  $G_N^2$  ? The equation for  $F_N^1$  depending on  $F_N^2$  also writes

$$\begin{aligned}\partial_t F_N^1(x) &= \frac{1}{2} \Delta F_N^1(x) - \operatorname{div}_x (b(x, F_N^1) F_N^1(x)) \\ &+ \operatorname{div}_x \left( \frac{1}{N} \int_{\mathbb{T}^d} \nabla W(x-y) (NG_N^2)(x, y) dy \right).\end{aligned}$$



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Assume that  $G_N^3 = O\left(\frac{1}{N^2}\right)$ , then the equation for  $F_N^2$  is

$$\begin{aligned}\partial_t F_N^2(x_1, x_2) &= \frac{1}{2} \Delta F_N^2(x_1, x_2) - \sum_{1 \leq i \neq j \leq 2} \operatorname{div}_{x_i} \left\{ -\frac{1}{N} \nabla W(x_i - x_j) F_N^1(x_i) F_N^1(x_j) \right. \\ &\quad + \frac{N-1}{N} b(x_i, F_N^1) F_N^1(x_i) F_N^1(x_j) + 3 \frac{N-1}{N} b(x_i, F_N^1) F_N^2(x_i, x_j) \\ &\quad - 3 \frac{N-1}{N} \int_{\mathbb{T}^d} \nabla W(x - x_i) F_N^2(x_i, x) dx F_N^1(x_j) \\ &\quad \left. - 3 \frac{N-1}{N} \int_{\mathbb{T}^d} \nabla W(x - x_i) F_N^2(x, x_j) dx F_N^1(x_1) \right\} + O\left(\frac{1}{N^2}\right).\end{aligned}$$

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Assume that  $G_N^3 = O\left(\frac{1}{N^2}\right)$ , then the equation for  $F_N^2$  is

$$\begin{aligned}\partial_t F_N^2(x_1, x_2) &= \frac{1}{2} \Delta F_N^2(x_1, x_2) - \sum_{1 \leq i \neq j \leq 2} \operatorname{div}_{x_i} \left\{ -\frac{1}{N} \nabla W(x_i - x_j) F_N^1(x_i) F_N^1(x_j) \right. \\ &\quad + \frac{N-1}{N} b(x_i, F_N^1) F_N^1(x_i) F_N^1(x_j) + 3 \frac{N-1}{N} b(x_i, F_N^1) F_N^2(x_i, x_j) \\ &\quad - 3 \frac{N-1}{N} \int_{\mathbb{T}^d} \nabla W(x - x_i) F_N^2(x_i, x) dx F_N^1(x_j) \\ &\quad \left. - 3 \frac{N-1}{N} \int_{\mathbb{T}^d} \nabla W(x - x_i) F_N^2(x, x_j) dx F_N^1(x_1) \right\} + O\left(\frac{1}{N^2}\right).\end{aligned}$$

Since  $G_N^2 = F_N^2 - (F_N^1)^{\otimes 2} \rightarrow$  closed form for the evolution of  $F_N^1$  and  $G_N^2$ . Initial data

1.  $G_N^2|_{t=0} = 0$ ;
2.  $F_N^1|_{t=0} = \mu_0$ .

## Controlling the correlations

Expect the contribution of  $G_N^2$  to be of order  $O\left(\frac{1}{N}\right)$ . With this contribution: correction to this mean-field limit, provided that  $G_N^3 = O\left(\frac{1}{N^2}\right)$ . And so on...

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Our work: in some weak sense and uniformly in time

$$G_N^{m+1} = O\left(\frac{1}{N^m}\right)$$

for all  $m \geq 1$ .

## A brief reminder on cumulants

(Joint) cumulants of  $(Z_1, \dots, Z_n)$  measure the **interactions** between the variables:  
for

$$K(t_1, \dots, t_n) = \log \mathbb{E} \left[ e^{\sum_{j=1}^n t_j Z_j} \right],$$

$$\kappa^n [Z_1, \dots, Z_n] = \frac{d^n}{dt_1 \dots dt_n} K(t_1, \dots, t_n) \Big|_{t_1 = \dots = t_n = 0}.$$

We write

$$\kappa^m(X) = \kappa^m(X, \dots, X).$$

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We write

$$\kappa^m(X) = \kappa^m(X, \dots, X).$$

Recall in particular that for all  $X \in L^4(\Omega)$ ,

$$\kappa^2(X) = \text{Var}(X), \quad \kappa^3(X) = \mathbb{E} \left[ (X - \mathbb{E}[X])^3 \right].$$

But of course it is not always that easy

$$\kappa^4(X) = \mathbb{E} \left[ (X - \mathbb{E}[X])^4 \right] - 3\text{Var}(X)^2.$$

# Main result

Our main result is the following:

## Theorem (B.-Duerinckx 2022<sup>+</sup>)

*Assume that  $b$  is given by a smooth,  $H$ -stable potential  $W$ , and that  $\Phi : \mathcal{P}(\mathbb{T}^d) \rightarrow \mathbb{R}$  is smooth. Then, for all  $m \geq 1$ , there exists a constant  $C > 0$  such that, for any  $\mu_0 \in \mathcal{P}(\mathbb{T}^d)$ ,*

$$\sup_{t \geq 0} \kappa^{m+1} \left[ \Phi(\mu_t^N) \right] \leq \frac{C}{N^m}.$$

- ⚡  $\Phi$  smooth in the sense of linear derivatives w.r.t. the measure.
- ⚡ Explicit dependency of  $C$  in the derivatives of  $\Phi$ .



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- ⚡  $\Phi$  smooth in the sense of linear derivatives w.r.t. the measure.
- ⚡ Explicit dependency of  $C$  in the derivatives of  $\Phi$ .

Possible to relate  $\kappa^{m+1}[\Phi(\mu_t^N)]$  to the norm of  $G_N^{m+1}$  when  $\Phi(\mu) = \int_{\mathbb{T}^d} \varphi(x) \mu(dx)$  with  $\varphi$  smooth.

## The sources of randomness

Recall  $\mu_t^N = \frac{1}{N} \sum_{i=1}^N \delta_{Y_t^{i,N}}$  for all  $t \geq 0$ . Let  $\Phi : \mathcal{P}(\mathbb{T}^d) \rightarrow \mathbb{R}$ . Weak formulation of the result for  $G_N^2$ :

$$\text{Var}[\Phi(\mu_t^N)] = O\left(\frac{1}{N}\right),$$

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Two sources of randomness, treated separately:

- Brownian motions;
- initial distributions.

$\mathbb{E}$  for the global randomness,  $\mathbb{E}_o$  for the one related to the initial data,  $\mathbb{E}_B$  for the one related to the Brownian motions. And so on, we write  $\text{Var}$ ,  $\text{Var}_o$ ,  $\text{Var}_B$ ,  $\kappa$ ,  $\kappa_o$ ,  $\kappa_B \dots$

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Splitting between those two sources:

$$\text{Var}[\Phi(\mu_t^N)] = \text{Var}_o[\mathbb{E}_B[\Phi(\mu_t^N)]] + \mathbb{E}_o[\text{Var}_B(\Phi(\mu_t^N))].$$

We will prove

- ☆  $\mathbb{E}_B[\Phi(\mu_t^N)] = \Phi(m(t, \mu_0^N)) + O\left(\frac{1}{N}\right)$ ;
- ☆  $\text{Var}_B(\Phi(\mu_t^N)) = O\left(\frac{1}{N}\right)$ .

## Our tools

Specific tools for each type of randomness.

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In both cases, ergodic estimates to obtain the uniform control in time.



## Linear functional derivatives

Let  $F : \mathcal{P}(\mathbb{T}^d) \rightarrow \mathbb{R}$ . We say that  $F$  is continuously differentiable if there exists a continuous function  $\frac{\delta F}{\delta m} : \mathcal{P}(\mathbb{T}^d) \times \mathbb{T}^d \rightarrow \mathbb{R}$  such that, for any  $\mu, \mu' \in \mathcal{P}(\mathbb{T}^d)$ ,

$$F(\mu) - F(\mu') = \int_0^1 \int_{\mathbb{T}^d} \frac{\delta F}{\delta m}(s\mu + (1-s)\mu', y)(\mu - \mu')(dy) ds.$$

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Wasserstein derivative: for  $y \in \mathbb{T}^d$ ,  $\mu \in \mathcal{P}(\mathbb{T}^d)$ ,

$$\partial_\mu F(\mu)(y) = \partial_y \frac{\delta F}{\delta m}(\mu, y).$$

## Glauber calculus

Let  $\gamma : (\mathbb{T}^d)^N \rightarrow \mathbb{R}$ . Glauber derivative with respect to  $Y_0^{1,N}$ :

$$D_{\circ}^1[\gamma(Y_0^{1,N}, \dots, Y_0^{N,N})] = \gamma(Y_0^{1,N}, \dots, Y_0^{N,N}) - \int_{\mathbb{T}^d} \gamma(z, \dots, Y_0^{N,N}) \mu_0(dz)$$

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$\implies$  measure the sensitivity of  $\gamma$  with respect to  $Y_0^{1,N}$ .

For any  $\psi : \mathcal{P}(\mathbb{T}^d) \rightarrow \mathbb{R}$  admitting linear derivative, any  $j \in [N]$ ,

$$\begin{aligned} D_0^j[\psi(\mu_0^N)] &= \frac{1}{N} \int_0^1 \frac{\delta\psi}{\delta m} \left( \frac{1}{N} \sum_{i \neq j} \delta_{Y_0^{i,N}} + \frac{s}{N} \delta_{Y_0^{j,N}} + \frac{1-s}{N} \delta_z, Y_0^{j,N} \right) \mu_0(dz) ds \\ &\quad - \frac{1}{N} \int_0^1 \frac{\delta\psi}{\delta m} \left( \frac{1}{N} \sum_{i \neq j} \delta_{Y_0^{i,N}} + \frac{s}{N} \delta_{Y_0^{j,N}} + \frac{1-s}{N} \delta_z, z \right) \mu_0(dz) ds \end{aligned}$$

$\implies D_0^j[\psi(\mu_0^N)] = O\left(\frac{1}{N}\right)$  provided good control of  $\frac{\delta\psi}{\delta m}$ .

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Efron-Stein's inequality:

$$\text{Var}_\circ[Y] \leq \mathbb{E}^\circ \left[ \sum_{j=1}^N |D_0^j[Y]|^2 \right].$$

Similar *Poincaré inequality* for higher-order cumulants.

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## The master equation

For any  $\Phi : \mathcal{P}(\mathbb{T}^d) \rightarrow \mathbb{R}$ , write  $\mathcal{U}_\Phi(t, \mu) = \Phi(m(t, \mu))$  for  $t \geq 0$ ,  $\mu \in \mathcal{P}(\mathbb{T}^d)$ . Then, from Buckdahn-Li-Peng-Rainer (2017),  $\mathcal{U}_\Phi$  satisfies the **master equation**

$$\begin{cases} \partial_t \mathcal{U}_\Phi(t, \mu) &= \int_{\mathbb{T}^d} \left[ \sum_{i=1}^d \partial_{x_i} \frac{\delta \mathcal{U}_\Phi}{\delta m}(t, \mu, x) b_i(x, \mu) \right. \\ &\quad \left. + \frac{1}{2} \sum_{i,j=1}^d \partial_{x_i x_j}^2 \frac{\delta \mathcal{U}_\Phi}{\delta m}(t, \mu, x) \right] \mu(dx) \quad t \geq 0, \\ \mathcal{U}_\Phi(0, \mu) &= \Phi(\mu) \end{cases}$$



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→ expand  $m(t, \mu)$  along the dynamics. From Chassagneux-Szpruch-Tse (2019), we have

$$\mathbb{E}_B[\Phi(\mu_i^N)] = \mathcal{U}_\Phi(t, \mu_0^N) + \frac{1}{2N} \int_0^t \int_{\mathbb{T}^d} \mathbb{E}_B \left[ \text{Tr}[\partial_\mu^2 \mathcal{U}_\Phi(t-s, \mu_s^N, v, v)] \mu_s^N(dv) \right] ds,$$

where  $\partial_\mu \mathcal{U}_\Phi(t-s, \mu, y) = \partial_y \frac{\delta \mathcal{U}}{\delta m}(t-s, \mu, y)$ .

## Pushing the expansion further

Set, for  $0 \leq s \leq t$ ,  $\mu \in \mathcal{P}(\mathbb{T}^d)$ ,

$$\Phi^{(1)}((t, s), \mu) = \int_{\mathbb{T}^d} \text{Tr} \left[ \partial_{\mu}^2 \mathcal{U}_{\Phi}(t - s, \mu, y, y) \right] \mu(dy),$$

and then set, for  $0 \leq u \leq s \leq t$ ,

$$\mathcal{U}_{\Phi}^{(1)}((t, s, u), \mu) = \Phi^{(1)}((t, s), m(s - u, \mu)).$$

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→ use  $\mathcal{U}_{\Phi}^{(1)}$  to push the expansion.

$$\begin{aligned} \mathbb{E}_B[\Phi(\mu_t^N)] &= \mathcal{U}_{\Phi}(t, \mu_0^N) + \frac{1}{2N} \int_0^t \mathcal{U}_{\Phi}^{(1)}((t, s, 0), \mu_0^N) ds \\ &\quad + \frac{1}{4N^2} \int_0^t \int_0^s \int_{\mathbb{T}^d} \text{Tr} \left[ \partial_{\mu}^2 \mathcal{U}_{\Phi}^{(1)}((t, s, u), \mu_u^N, y, y) \right] \mu_u^N(dy) du ds. \end{aligned}$$

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**Explicit formulas** relating  $\partial_{\mu}^2 \mathcal{U}_{\Phi}^{(1)}$  with Wasserstein derivatives of  $\Phi$  evaluated at solutions of linearized parabolic equations. In particular, using **ergodic estimates** for those solutions:

$$\sup_{\mu \in \mathcal{P}(\mathbb{T}^d)} \int_{\mathbb{T}^d} \int_0^t \int_0^s \text{Tr} \left[ \partial_{\mu}^2 \mathcal{U}_{\Phi}^{(1)}((t, s, u), \mu, y, y) \right] \mu(dy) du ds = O(1).$$

## Treating the Brownian cumulants

$$\begin{aligned}\text{Var}_B[\Phi(\mu_t^N)] &= \mathbb{E}_B[\Phi(\mu_t^N)^2] - \mathbb{E}_B[\Phi(\mu_t^N)]^2 \\ &= \mathcal{U}_{\Phi^2}(t, \mu_0^N) + \frac{1}{2N} \int_0^t \mathcal{U}_{\Phi^2}^{(1)}((t, s, 0), \mu_0^N) ds \\ &\quad - \left( \mathcal{U}_{\Phi}(t, \mu_0^N) + \frac{1}{2N} \int_0^t \mathcal{U}_{\Phi}^{(1)}((t, s, 0), \mu_0^N) ds \right)^2 + O\left(\frac{1}{N^2}\right)\end{aligned}$$

and  $\mathcal{U}_{\Phi^2}(t, \mu_0^N) = \Phi^2(m(t, \mu_0^N)) = \mathcal{U}_{\Phi}(m(t, \mu_0^N))^2$  so  $\text{Var}_B[\Phi(\mu_t^N)] = O\left(\frac{1}{N}\right)$ .

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We can do much more ! Identifying precisely the  $O\left(\frac{1}{N}\right)$  term: since

$$\partial_{\mu}^2 f^2(\nu)(x, x) = 2f(\nu)(\partial_{\mu}^2 f(\nu)(x, x)) + 2(\partial_{\mu} f(\nu)(x))^2$$

$$\begin{aligned}\mathcal{U}_{\Phi^2}^{(1)}((t, s, 0), \mu_0^N) &= 2\mathcal{U}_{\Phi}(t, \mu_0^N)\mathcal{U}_{\Phi}^{(1)}((t, s, 0), \mu_0^N) \\ &\quad + \int_{\mathbb{T}^d} |\partial_{\mu}\mathcal{U}_{\Phi}(t-s, m(s, \mu_0^N))(y)|^2 m(s, \mu_0^N)(dy).\end{aligned}$$

## Treating the Brownian cumulants

$$\begin{aligned}\text{Var}_B[\Phi(\mu_t^N)] &= \mathbb{E}_B[\Phi(\mu_t^N)^2] - \mathbb{E}_B[\Phi(\mu_t^N)]^2 \\ &= \mathcal{U}_{\Phi^2}(t, \mu_0^N) + \frac{1}{2N} \int_0^t \mathcal{U}_{\Phi^2}^{(1)}((t, s, 0), \mu_0^N) ds \\ &\quad - \left( \mathcal{U}_{\Phi}(t, \mu_0^N) + \frac{1}{2N} \int_0^t \mathcal{U}_{\Phi}^{(1)}((t, s, 0), \mu_0^N) ds \right)^2 + O\left(\frac{1}{N^2}\right)\end{aligned}$$

and  $\mathcal{U}_{\Phi^2}(t, \mu_0^N) = \Phi^2(m(t, \mu_0^N)) = \mathcal{U}_{\Phi}(m(t, \mu_0^N))^2$  so  $\text{Var}_B[\Phi(\mu_t^N)] = O\left(\frac{1}{N}\right)$ .

We can do much more ! Identifying precisely the  $O\left(\frac{1}{N}\right)$  term: since

$$\partial_{\mu}^2 f^2(\nu)(x, x) = 2f(\nu)(\partial_{\mu}^2 f(\nu)(x, x)) + 2(\partial_{\mu} f(\nu)(x))^2$$

$$\begin{aligned}\mathcal{U}_{\Phi^2}^{(1)}((t, s, 0), \mu_0^N) &= 2\mathcal{U}_{\Phi}(t, \mu_0^N)\mathcal{U}_{\Phi}^{(1)}((t, s, 0), \mu_0^N) \\ &\quad + \int_{\mathbb{T}^d} |\partial_{\mu}\mathcal{U}_{\Phi}(t-s, m(s, \mu_0^N))(y)|^2 m(s, \mu_0^N)(dy).\end{aligned}$$

Hence,

$$\text{Var}_B[\Phi(\mu_t^N)] = \frac{1}{N} \int_0^t \int_{\mathbb{T}^d} |\partial_{\mu}\mathcal{U}_{\Phi}(t-s, m(s, \mu_0^N), y)|^2 m(s, \mu_0^N)(dy) ds + O\left(\frac{1}{N^2}\right).$$

Can apply Glauber calculus to this leading term !

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Thank you for your attention !